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Lim, Hyang-Tag; Togan, Emre; Kroner, Martin; Miguel-Sanchez, Javier; Imamoglu, Atac

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Electrically tunable artificial gauge potential for polaritons

Hyang-Tag Lim^{1,*}, Emre Togan^{1,*}, Martin Kroner¹, Javier Miguel-Sanchez¹ & Atac Imamoglu¹

Neutral particles subject to artificial gauge potentials can behave as charged particles in magnetic fields. This fascinating premise has led to demonstrations of one-way waveguides, topologically protected edge states and Landau levels for photons. In ultracold neutral atoms, effective gauge fields have allowed the emulation of matter under strong magnetic fields leading to realization of Harper-Hofstadter and Haldane models. Here we show that application of perpendicular electric and magnetic fields effects a tunable artificial gauge potential for two-dimensional microcavity exciton polaritons. For verification, we perform interferometric measurements of the associated phase accumulated during coherent polariton transport. Since the gauge potential originates from the magnetoelectric Stark effect, it can be realized for photons strongly coupled to excitations in any polarizable medium. Together with strong polariton-polariton interactions and engineered polariton lattices, artificial gauge fields could play a key role in investigation of non-equilibrium dynamics of strongly correlated photons.

¹Institute of Quantum Electronics, ETH Zurich, CH-8093 Zurich, Switzerland. * These authors contributed equally to this work. Correspondence and requests for materials should be addressed to E.T. (email: togane@phys.ethz.ch) or to A.I. (email: imamoglu@phys.ethz.ch).

Synthesis of artificial gauge fields for photons have been demonstrated in a number of different optical systems. In most cases, the implementation is achieved through the design of the optical system^{1–4}, leaving little or no room for fast control of the magnitude of the effected gauge field after sample fabrication is completed. For many applications on the other hand, it is essential to be able to tune or adjust the strength of the gauge field during the experiment^{5–7}; this is particularly the case for nanophotonic structures⁸ where fast local control of the gauge field strength may open up new possibilities for investigation of many-body physics of light⁹.

Cavity-polaritons are hybrid light-matter quasi-particles arising from non-perturbative coupling between quantum well (QW) excitons and cavity photons. A magnetic field $\mathbf{B}=B_z\hat{z}$ applied along the growth direction influences polaritons through their excitonic nature and leads to a diamagnetic shift, Zeeman splitting of circularly polarized modes and enhancement of exciton-photon coupling strength¹⁰. For excitons with a non-zero momentum, the applied B_z also induces an electric dipole moment^{11–14}. In a classical picture, this dipole moment is due to the Lorentz force that creates an effective electric field \mathbf{E}_{eff} (ref. 15) that is proportional to $\mathbf{k} \times \mathbf{B}$. For excitons with small momentum and polarizability α , this \mathbf{E}_{eff} causes an induced dipole moment ($\mathbf{d} \propto \alpha \mathbf{k} \times \mathbf{B}$). An additional external electric field \mathbf{E}_{ext} in the QW plane then alters the dispersion of excitons as it leads to energy changes proportional to k : $-\mathbf{d} \cdot \mathbf{E}_{\text{ext}} \propto -\alpha(\mathbf{k} \times \mathbf{B}) \cdot \mathbf{E}_{\text{ext}}$, so that the energy minimum is no longer at $|\mathbf{k}|=0$ as would be the case for a free particle with an effective mass but at finite \mathbf{k} . This simple modification of the dispersion relation is equivalent to an effective gauge potential \mathbf{A}_{eff} for excitons; due to their partly excitonic character, the Hamiltonian describing the dynamics of polaritons also contains an effective gauge potential term proportional to \mathbf{A}_{eff} .

In the realization we describe here, based on the effective gauge potential \mathbf{A}_{eff} , the strength and direction of the effected gauge potential is controlled electrically. Moreover, since our scheme relies on the magnetoelectric Stark effect, time reversal symmetry^{16–18} of the optical excitations is broken. This realization, in combination with strong polariton-polariton interactions^{19–21}, has the potential to open up new possibilities of investigating many-body physics of light.

Results

Demonstration of the effective electric field \mathbf{E}_{eff} . The cavity polariton sample we use is illustrated in Fig. 1a. Our demonstration relies on the fact that it is possible to excite polaritons with a well-defined in-plane wavevector \mathbf{k} by appropriately choosing the angle and energy of the excitation beam. As illustrated in Fig. 1a, we choose $\mathbf{k}=k_y\hat{y}$. Two metal gates deposited 30 μm apart allow us to apply an electric field in the x direction such that for polaritons propagating with a wavevector along the y direction, $\mathbf{E}_{\text{ext}}=E_{\text{ext}}\hat{x}$ will add or cancel $\mathbf{E}_{\text{eff}}=E_{\text{eff}}\hat{x}$ due to the Lorentz force. By recording changes in the transition energy of polaritons as a function of E_{ext} , we can determine the strength of E_{eff} . In our experiments, we only address the lower polariton branch.

Changes in the reflected intensity of a laser beam probing polaritons at $B_z=0$ T with $k_y^+=2.7\ \mu\text{m}^{-1}$ as a function of E_{ext} is shown in Fig. 1d. For each E_{ext} the reflected intensity shows a dip at an energy corresponding to the polariton resonance (Fig. 1e). The spectral centre of the dip shifts to lower energies with E_{ext} in a way that is well described by a second-order polynomial. The expected behaviour for a neutral polarizable quasi-particle (that is, exciton or polariton) that is subject to E_{ext} , would be to have a dc-Stark shift equal to $-\alpha|E_{\text{ext}}|^2$ (ref. 22); here

the coefficient α is the quasi-particle polarizability. Therefore, the electric field that yields the maximum lower-polariton energy identifies $E_{\text{ext}}=\tilde{E}_{\text{ext}}$ that exactly cancels any internal or effective electric fields E_{eff} ($\tilde{E}_{\text{ext}}=-E_{\text{eff}}$). To quantify E_{eff} due to the Lorentz force described above, we extract the difference of \tilde{E}_{ext} at a fixed B_z for polaritons excited with $k_y^+=2.7\ \mu\text{m}^{-1}$ and $k_y^-=-2.9\ \mu\text{m}^{-1}$. This approach also allows us to exclude influences of built-in electric fields. At $B_z=0$ T, we find that the difference of V_G values that correspond to the maximum energy with k_y^\pm excitation is 0.02 V, indicating that E_{eff} is negligible at 0 T (Fig. 1f).

In stark contrast, for $B_z=5$ T, the energy shift of polaritons with E_{ext} displays a significant difference between \tilde{E}_{ext} for k_y^+ and k_y^- , as illustrated in Fig. 2a. The magnetic field dependence of the difference of \tilde{E}_{ext} for k_y^\pm , Fig. 2b, shows a behaviour that is well described by a linear increase with B_z , demonstrating E_{eff} due to the Lorentz force.

B_z -induced changes to polaritons. With increasing B_z , the polarizability decreases and the polariton energy at $E_{\text{ext}}=0$ increases, as illustrated in Fig. 2c. The polariton energy at $E_{\text{ext}}=0$ and the polarizability can be extracted from a fit to a second-order polynomial: $\epsilon_0 - dE_{\text{ext}} - \alpha E_{\text{ext}}^2$. The first-order coefficient of the polynomial in turn, yields the induced dipole moment of the polaritons (excitons). The difference of the first-order coefficient for $\pm k_y$ excitation beams (twice the induced dipole moment) is shown in Fig. 2d: we observe that the induced dipole moment first increases with B_z and then starts to slowly decrease for $B_z \geq 4$ T. We expect that for the high magnetic field regime, the induced dipole moment decreases as $|d| \sim 1/B_z$ (ref. 12). We find very good agreement with a theoretical model of the polaritons (shown as solid lines in Fig. 2) that allows us to identify the physical mechanism underlying the overall B_z and E_{ext} dependence (see Supplementary Note 4 for detailed information about the model). The change in the lower polariton energy is due to an interplay between a diamagnetic blue-shift of the exciton transition energy and a red-shift due to an increase in the cavity-exciton coupling strength¹⁰. Concurrently, the polarizability decreases with B_z , as would be expected in a classical picture as the size of the polarizable particle decreases. We emphasize that due to the change in the exciton energy and cavity-exciton coupling strength the exciton content of the polaritons changes as B_z is varied. We also find that the energy shift of polaritons with E_{ext} is influenced by the decrease in the electron-hole overlap due to E_{ext} -induced polarization, which leads to a reduction of the cavity-exciton coupling strength. In the E_{ext} range that we probe the decrease in the cavity-exciton coupling leads to an effective smaller polarizability for polaritons as compared to excitons, due to a reduction on the excitonic character of the polaritons. Moreover, the model also captures the presence of E_{eff} due to the centre of mass motion of polaritons and confirms that even with changes in exciton content with B_z , its expected dependence on B_z is still linear. Finally, we emphasize that while E_{eff} arises from exciton physics, the strong coupling between excitons and photons has to be taken into account to obtain a quantitative agreement between the model and our experiments.

Demonstration of a gauge potential for polaritons. Remarkably, our findings demonstrate that polariton energy under perpendicular \mathbf{B} and \mathbf{E}_{ext} depends not only on the magnitude of \mathbf{k} but also on its direction: this non-reciprocal flow of light is an indication of the presence of a gauge field for polaritons. The dispersion relation for free excitons propagating in the QW plane with wavevector $k_y\hat{y}$ in the presence of an electric field

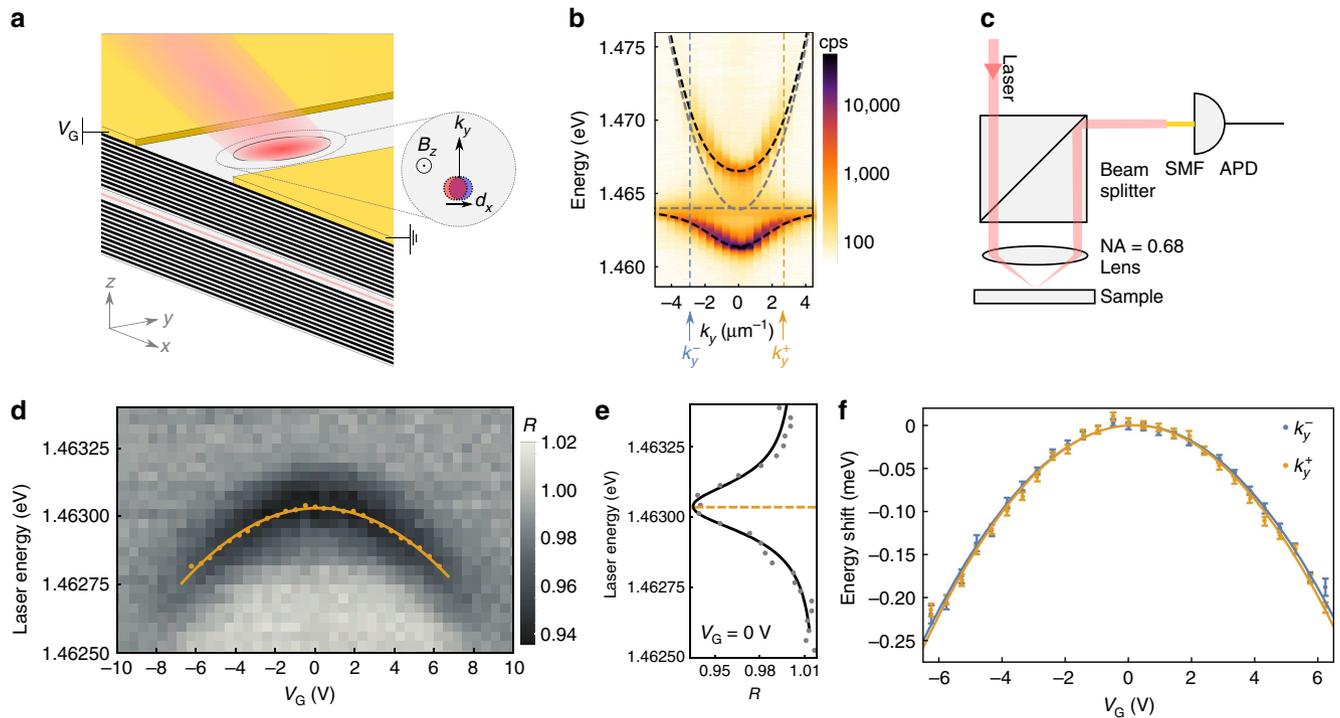


Figure 1 | Description of the sample and characterization experiments at $B_z = 0$ T. (a) The sample, held at 4 K in a helium bath cryostat, contains three $\text{In}_{0.04}\text{Ga}_{0.96}\text{As}$ quantum wells located at an antinode of a cavity formed by two DBRs. An electric potential V_G applied to the metal gates deposited on the sample surface creates an electric field in the x direction (see ‘Methods’ section). Inset: with a magnetic field in the z direction polaritons excited with in-plane wavevector k_y exhibit a dipole moment d_x . (b) k_y -resolved PL spectra at $B_z = 0$ T. Grey dashed lines show the extracted bare cavity and exciton dispersions, black lines show polariton dispersions (see Supplementary Note 3 for parameters), orange and blue lines correspond to $k_y = k_y^+$ and $k_y = k_y^-$, respectively. cps stands for counts per second. (c) Experimental set-up used in reflection measurements. Reflection of a linearly polarized laser beam (< 100 pW) with in-plane wavevector k_y is coupled to a single-mode fibre (SMF) and its intensity is detected with an avalanche-photodiode (APD). (d) Changes in the reflected intensity, R , of a laser beam exciting polaritons with in-plane wavevector k_y^+ as a function of laser energy and V_G . Yellow dots indicate the extracted polariton resonance energies. Note that the polariton signal is missing ($|V_G| > 8$ V) due to the ionization of the exciton. (e) Line cut of the reflection spectrum at $V_G = 0$ V. Grey points are the measured reflection data, black solid line shows the fitted lineshape (see Supplementary Note 3). Dashed yellow line indicates the extracted polariton resonance energy. (f) Change of polariton energy from $V_G = 0$ V, as a function of V_G for polaritons excited with k_y^+ (yellow) and k_y^- (blue) at $B_z = 0$ T, solid lines are fits to second-order polynomials (blue: $2.1 \times 10^{-6} V_G - 5.6 \times 10^{-6} V_G^2$, yellow: $1.9 \times 10^{-6} V_G - 5.8 \times 10^{-6} V_G^2$). Error bars represent 1 s.d. of the polariton resonance energy extracted from a fit of the data to a reflection lineshape (see Supplementary Note 3).

$$\mathbf{E} = \left(\frac{\hbar}{M} k_y B_z + E_{\text{ext}} \right) \hat{\mathbf{x}} \text{ is:}$$

$$\epsilon_{\text{exc}}(k_y) = \epsilon_{\text{exc}}(0) + \frac{\hbar^2 k_y^2}{2M} - \alpha \left(\frac{\hbar}{M} k_y B_z + E_{\text{ext}} \right)^2 = \epsilon'_{\text{exc}} + \frac{(\hbar k_y - qA_{\text{eff}})^2}{2M'}, \quad (1)$$

where M is the total exciton mass, and $\epsilon_{\text{exc}}(0)$ is the exciton energy at $k_y = 0$ and $E_{\text{ext}} = 0$. $qA_{\text{eff}} = \frac{M'}{M} 2\alpha B_z E_{\text{ext}}$ is an effective vector potential for excitons, with $\epsilon'_{\text{exc}} = \epsilon_{\text{exc}}(0) - \alpha E_{\text{ext}}^2 - \frac{(qA_{\text{eff}})^2}{2M'}$ and $M' = M \left(1 - \frac{2\alpha B_z^2}{M} \right)^{-1}$. The strong coupling to the cavity further changes the dispersion relation, but polaritons in a narrow energy-momentum range can still be treated as free particles under the influence of an effective vector potential for polaritons, qA , whose value depends on the detuning between the excitons and polaritons as well as their effective masses (see Supplementary Note 4).

For our experiments, both B_z and E_{ext} are uniform over the area of our experiment yielding a constant gauge potential qA . Since physical observables are gauge invariant, we might expect that a constant gauge potential would not have any observable experimental signatures, and our measurements to be independent of qA . In practice, our experiments involve propagation of photons between two regions with

different constant gauge potentials; when calculating the magnitude of qA above, we fixed a gauge by choosing the constant gauge potential for photons outside the cavity to be $qA_{\text{out}} = 0$. The difference $qA - qA_{\text{out}} \neq 0$ is gauge invariant and stems from an effective (sheet) magnetic field at the interface between the cavity and vacuum. The experimental signatures of qA , such as shift of the energy minimum of polariton dispersion away from $k = 0$, are observable in our experiments due to this effective magnetic field that alters the kinetic momentum as photons/polaritons pass through the interface^{23–25}.

To demonstrate the validity of the description of photon/polariton dynamics as determined by a gauge potential, we perform an interference experiment where we measure the phase accumulated by polaritons propagating for a length l inside the sample relative to a fixed phase reference (see Fig. 3a and ‘Methods’ section for detailed information about the interference experiment). An interference image obtained at $E_{\text{ext}} = 0$ is shown in Fig. 3b. Changes in the interference pattern with y at two different values of E_{ext} are illustrated in Fig. 3c.

A plot of the phase change of polaritons with E_{ext} for different directions of excitation at $B_z = 6$ T is shown in Fig. 3d. The data show that at a fixed E_{ext} and at the same excitation frequency,

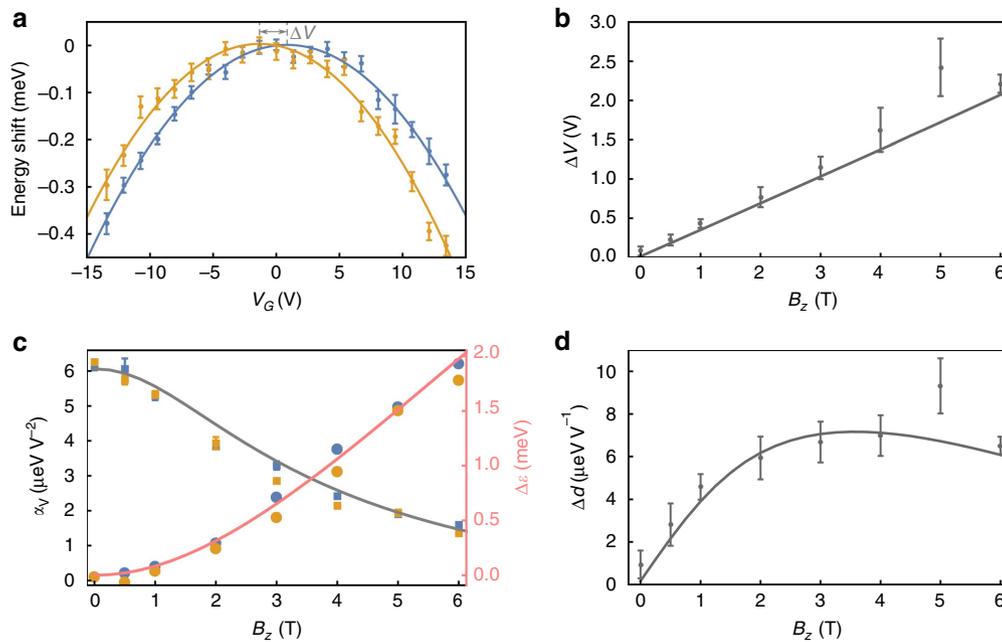


Figure 2 | Magnetic field dependence. (a) Shift of polariton energies with applied voltage for polaritons excited with k_y^+ (yellow) and k_y^- (blue) at $B_z = 5$ T; solid lines are fits to second-order polynomials ($-d_V V_G - \alpha_V V_G^2$; blue: $d_V = 3.0 \times 10^{-6} \text{ eV V}^{-1}$, $\alpha_V = 1.8 \times 10^{-6} \text{ eV V}^{-2}$, yellow: $d_V = -5.2 \times 10^{-6} \text{ eV V}^{-1}$, $\alpha_V = 1.9 \times 10^{-6} \text{ eV V}^{-2}$). Error bars represent 1 s.d. of the polariton resonance energy extracted from a fit of the data to a reflection lineshape (see Supplementary Note 3). Grey dashed lines indicate the voltage at which maximum energy occurs for the two parabolas, the grey arrow indicates the difference between them (ΔV). (b) Using data such as **a** at different B_z , difference between the extracted voltage (ΔV) at which maximum energy occurs for k_y^- and k_y^+ excitations. (c) Change in α_V (polarizability) with B_z (filled squares, left axis, and grey curve). Change of lower polariton resonant energy with B_z at $V_G = 0$ V (filled circles, right axis, and red curve) from the value at $B_z = 0$ T. Yellow and blue data points are data for k_y^+ and k_y^- , respectively. (d) Difference between extracted d_V values (electric dipole moment) for polaritons propagating with k_y^+ and k_y^- . For **b-d**, error bars are the estimated standard deviations of the mean for three repetitions of the experiment (some error bars are smaller than the markers), and the solid lines are results of a numerical calculation (see Supplementary Note 4).

the phase accumulated for polaritons propagating along $\pm \hat{y}$ directions differ, and that this phase difference depends on the sign of E_{ext} and B_z . The kinetic momentum of polaritons with a fixed energy is modified by the absence or presence of a constant vector potential so that the additional accumulated phase for travelling a distance l is given by $\frac{q}{\hbar} \int \mathbf{A} \cdot d\mathbf{l}$. At a fixed E_{ext} , the difference between the phase accumulated when exciting with k_y^\pm allows us to avoid phase changes due to the Stark effect and to directly obtain $\Delta\phi_{k_y^+} - \Delta\phi_{k_y^-} = 2\frac{q}{\hbar} Al$. The fact that $\Delta\phi_{k_y^+} - \Delta\phi_{k_y^-}$ depends both on the sign of the electric field as well as the magnetic field is another manifestation of the validity of the description with qA . Since the phase shift as a function of E_{ext} can be modelled with a second-order polynomial in E_{ext} , the dc-Stark effect is the dominant underlying mechanism for the phase shift. Figure 3d shows that polaritons acquire a phase shift of ~ 0.25 radians due to qA as they propagate an average distance $l = 9 \mu\text{m}$.

Discussion

The parameters of our experiment already ensure significant phase accumulation over the lattice constant ($\geq 2.4 \mu\text{m}$) of state-of-the-art polariton lattices⁸ or average separation ($\sim 4 \mu\text{m}$) of coupled polariton microcavities²⁶. Choosing detunings that yield predominantly exciton-like lower polaritons will further increase qA . In addition, employing structures with longer polariton lifetimes²⁷ will allow the realization of polariton lattices with larger lattice constants. With these advances, it should be possible to design structures where the accumulated polariton phase due to qA in a unit cell could be on the order of π .

Together with strong photon–photon interactions, realization of tunable artificial gauge fields is key for ongoing research aimed at the observation of topological order in driven-dissipative photonic systems^{28,29}. Cavity-polaritons constitute particularly promising candidates in this endeavour since their partly excitonic nature enhances interactions^{19–21}, and in an external magnetic field interplay of exciton Zeeman splitting and photonic spin-orbit coupling, can lead to opening of non-trivial energy gaps^{30,31}. Our work demonstrates that the same polariton system may be complemented by the realization of tunable gauge fields, if an in-plane electric-field gradient is introduced in the QW. While such gradients are manifest in most gated structures, creation of a constant artificial magnetic field over a region of several microns should be possible in specially designed structures. We envision that lattices of dipolaritons³² would provide a very promising avenue towards this endeavour: the use of an in-plane magnetic field in this case would effect a gauge potential that is substantially enhanced by the large dipolariton polarizability and strong external electric fields along the growth direction^{13,14}. The gradients of the latter in turn could be used to generate fluxes approaching magnetic flux quantum in a plaquette of area $\sim 5 \times 5 \mu\text{m}^2$. (Di)Polaritons that are subject to complex electric field distributions enabling large, tunable and inhomogeneous effective gauge field profiles³³ would constitute a novel feature in the exploration of topologically non-trivial non-equilibrium quantum systems.

Methods

Sample. Our sample consists of three layers of 9.6 nm-thick $\text{In}_{0.04}\text{Ga}_{0.96}\text{As}$ QWs sandwiched between two distributed Bragg reflectors (DBRs) formed by 20 (top)

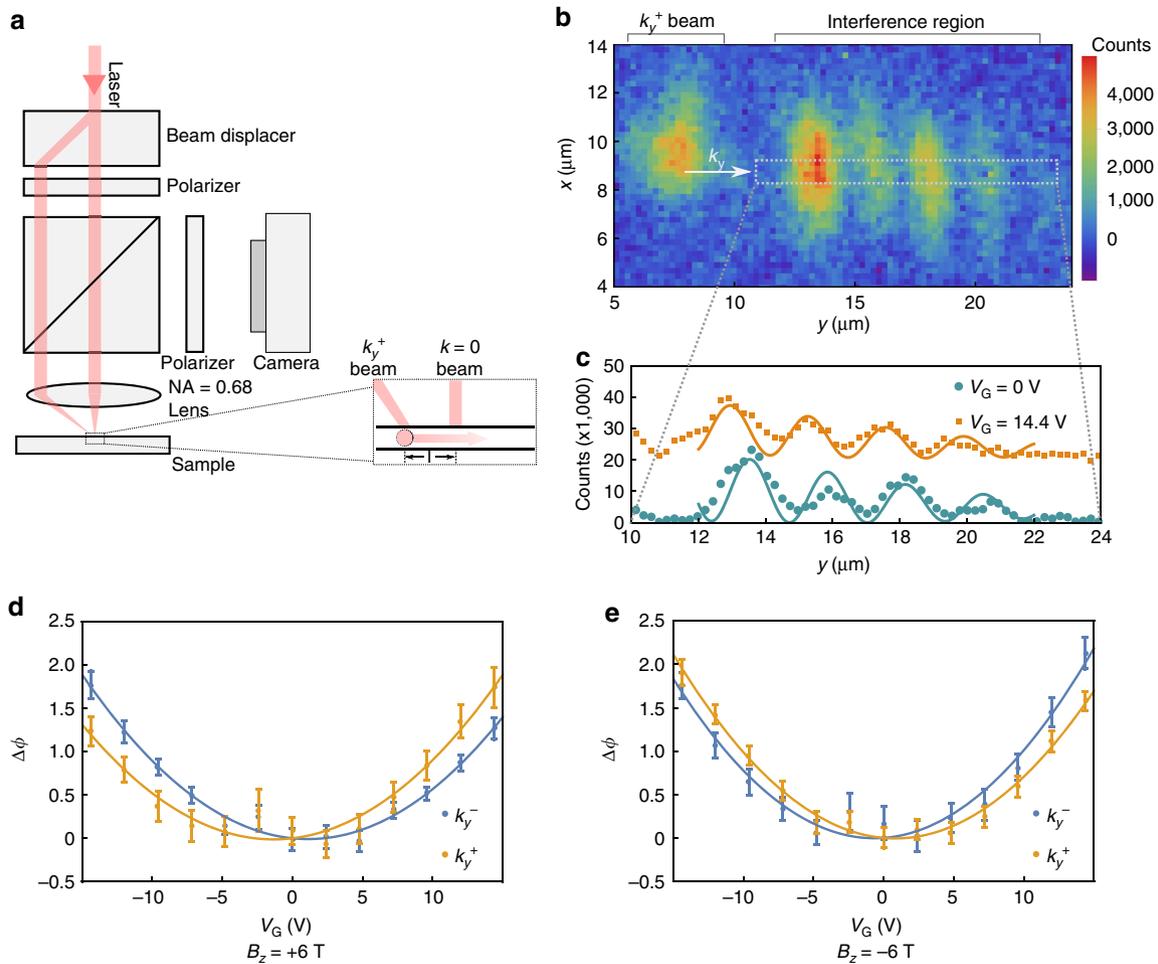


Figure 3 | Demonstration of effective gauge potential for propagating polaritons. (a) A single laser beam is split into two using a birefringent beam displacer, ensuring the relative phase between the two beams does not fluctuate during the experiment. Resulting two linearly polarized beams (total power 40 nW) are incident at different positions on the high numerical aperture (NA = 0.68) lens. One beam that is 1.2 mm off from the centre of the lens is incident on the sample with k_y^\pm . The second beam passes through the centre of the lens and thus has vanishing in-plane momentum ($k=0$ beam). The sample-lens distance is less than the focal length of the lens to ensure that the two beams are incident on the sample with their centres displaced by $l \sim 9 \mu\text{m}$. To detect polariton flow, we use an additional linear polarizer between the camera and sample that transmits light that is nearly orthogonally polarized to the incident beams. (b) An interference image obtained at $B_z = -6 \text{ T}$, $V_G = 0 \text{ V}$. Polaritons excited by the k_y^+ beam at 1.46467 eV propagate in the direction indicated by the white arrow. In the spatial region where photons emitted by propagating polaritons overlap with the reflected $k=0$ beam an interference pattern is observed. (c) In green, sum of the detected intensity of the interference pattern shown in b in the range $x = [8.5 - 9.4] \mu\text{m}$. In yellow, same as for green but shifted in intensity, for an image obtained at $V_G = 14.4 \text{ V}$. Solid lines are fits to the model described in ‘Methods’ section and show a difference in the extracted phase. (d) Change in the extracted phase at $B_z = 6 \text{ T}$ as a function of V_G for polaritons excited with k_y^+ (yellow) and k_y^- (blue). Solid lines show second-order polynomial fits to the phase change as a function of V_G (blue: $-0.016 V_G + 0.0072 V_G^2$, yellow: $0.019 V_G + 0.0071 V_G^2$). Error bars represent 1 s.d. in parameter estimation of the phase, when polariton intensity is fit to equation (2) (see ‘Methods’ section). (e) Same as d at $B_z = -6 \text{ T}$, blue: $0.011 V_G + 0.0089 V_G^2$, yellow: $-0.014 V_G + 0.0084 V_G^2$.

and 22 (bottom) pairs of $\lambda/4$ thick AlAs/GaAs layers. The GaAs spacer layer containing QWs between the top and bottom DBRs is 1λ thick. Spectral linewidths, measured at a position where the cavity mode is far red-detuned from the exciton mode, indicate that the cavity has a Q factor exceeding 10^4 (full-width at half-maximum linewidth 0.11 meV). The exciton emission spectrum at the same position has a full-width at half-maximum linewidth less than 1.2 meV.

Using electron beam lithography and lift-off techniques, we have defined rectangular metal (10 nm Ti, 210 nm Au) pads that are $700 \times 120 \mu\text{m}$. All of the experiments reported in this paper were carried out between two such pads that are separated by $30 \mu\text{m}$. These pads are wire bonded onto a chip carrier, and were driven by a high-voltage amplifier (Falco Systems WMA-300) to apply electric fields to the QW. The position on the sample is chosen such that for $|\mathbf{k}| \sim 0$ and $B_z = 0$ the detuning of the cavity mode energy from the exciton resonance is 0.09 meV, which is small compared to the exciton cavity coupling strength 5.2 meV.

In-plane electric field. As described in Fig. 1a, we apply an electric potential V_G between two gates to create E_{ext} . We find a good agreement between our data

and a numerical calculation for the polariton behaviour (see Supplementary Notes 1 and 4) with the electric field $E_{\text{ext}} = 0.84 V_G / (30 \mu\text{m})$. Since the two quantities are linearly related, we use E_{ext} and V_G interchangeably. For all experiments that involve a non-zero in-plane electric field, we acquire data in a pulse sequence (see Supplementary Note 2). This sequence is applied to avoid charge accumulation-induced variation in the actual electric field.

k-resolved photoluminescence spectra. To measure the k -resolved photoluminescence (PL) spectra, polaritons are created non-resonantly by a 780 nm continuous-wave diode laser. The incoming collimated laser beam passes through a high numerical aperture (NA) lens (NA = 0.68) so that the beam is focused on the sample plane. The PL signals from the sample also pass through the same lens. The back aperture of the high NA lens is the Fourier plane of emission. Emission around a particular in-plane momentum (k_x, k_y) is collected by coupling the light from a small spot in the Fourier plane into a single-mode fibre which is connected to a spectrometer. To obtain the polariton energy dispersion with respect to the in-plane wavevector k_y , we varied the position of the single-mode fibre, which

changes the collection spot on the Fourier plane. For each position that corresponds to $k_x = 0$ and a k_y value, we record the PL spectrum on the spectrometer. The incident 780 nm laser is blocked using a 800 nm long-pass filter at the entrance to the spectrometer.

Polariton interference measurement. We use the set-up depicted in Fig. 3a to create two beams that are incident with different in-plane momenta (k_y^\pm and $k = 0$) and are separated by l . The laser energy is chosen to be nearly resonant with the polariton modes with k_y^\pm . Owing to the steep polariton dispersion, at the same energy, the $k = 0$ beam is off resonance from the polariton modes; the photons in this beam are therefore reflected by the top mirror without being subject to qA . Photons in the k_y^\pm beam are converted into a propagating polariton cloud. As polaritons propagate towards the position of the $k = 0$ beam, they accumulate an additional phase due to qA . Propagating polaritons continue to emit photons out of the sample, and when these photons overlap with those from the $k = 0$ beam they interfere, allowing us to measure changes in the accumulated phase.

We model the light intensity in the region where the interference occurs as the sum of a Gaussian ($k = 0$ beam) and a plane wave (polaritons propagating with k_y):

$$\left| b_0 e^{-(y-y_0)^2/\sigma_y^2} + a_0 e^{i\phi} e^{-ik_y y} \right|^2 + b_1. \quad (2)$$

We independently determine b_0 , σ_y , and y_0 using images obtained at very high applied electric field, for which the polariton term is negligible ($a_0 \ll b_0$), and fix k_y . For each E_{ext} , we fit the detected polariton intensity to equation (2) to find values of a_0 , b_1 and ϕ .

Since our excitation beams have a finite size, they have a finite width in k_y ($0.45 \mu\text{m}^{-1}$), therefore, the exact k_y value of the polaritons that are excited is determined by the energy of the excitation beam. For a fixed excitation laser energy, changes in E_{ext} will shift the dispersion relation so that $k_y^\pm \Delta k_y$ is resonant with the excitation laser. These changes (due to qA or Stark effect) in the dispersion relation show up as phase changes in our experiments, $\Delta\phi = l\Delta k_y$, where l is the mean distance from the excitation spot to the interference region. Since we are interested in how the phase changes with E_{ext} and the absolute phase between the two beams is not determined, we set the phase at $E_{\text{ext}} = 0$ to 0 radian.

Data availability. The data that support the plots within this paper and other findings of this study are available from the corresponding author upon reasonable request.

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Author contributions

H.-T.L. and E.T. designed the experiments, carried out the measurements and analysed the results. J.M.-S. and E.T. fabricated the sample. M.K. helped with the experiments. A.I. supervised the project. H.-T.L., E.T. and A.I. wrote the manuscript.

Additional information

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