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# Some complementary estimates in the Dead Core problem

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Seminar für Angewandte Mathematik  
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## **Abstract**

This article is concerned with some estimates in the problem  $\Delta u = c^2 u^p$  in  $\Omega$ ,  $u = 1$  on  $\partial\Omega$ . The quantities of interest here are the critical value of  $c$  for which a dead core  $\Omega_0$  exists, the location of  $\Omega_0$  and the effectiveness factor.

**Keywords:** dead core, diffusion reaction

**Subject Classification:** 35J60

# 1 Introduction

This article is concerned with the problem

$$(1.1) \quad \begin{cases} \Delta u = c^2 u^p & \text{in } \Omega \subset \mathbb{R}^N, \\ u = 1 & \text{on } \partial\Omega, \end{cases}$$

with  $p \in (0, 1)$ . This precise form of the nonlinearity was chosen for simplicity only and in fact one can replace  $u^p$  by a more general function  $f(u)$  satisfying

$$f(0) = 0, \quad f'(s) \geq 0 \quad \text{and} \quad \int_0^1 \frac{ds}{f(s)} < \infty.$$

It was shown in [3] that for sufficiently large  $c$  a “dead core”  $\Omega_0$  develops in  $\Omega$ , i.e. a region where  $u \equiv 0$ . Problem (1.1) and various generalizations have been studied since by many different authors. There is no attempt made here of giving a survey of the literature on this problem and only the papers directly related to the results of this paper will be cited.

The quantities of interest for which estimates will be derived are:

- the critical value  $c_0$  of  $c$  above which a dead core will exist,
- the location of the dead core,
- the “effectiveness factor”  $\eta := \frac{\int_{\Omega} u^p dx}{|\Omega|}$ .

The estimates to be derived complement the corresponding ones given in [3].

## 2 Estimates derived from optimal supersolutions

In the following an essential assumption made is that the mean curvature of  $\partial\Omega$  be nonnegative everywhere. Most of the estimates will in general no longer hold if this assumption is dropped, as counterexamples show. Hence, this is not merely a purely technical assumption.

We now construct an optimal supersolution by combining the one-dimensional version of (1.1) with a suitable linear problem defined on  $\Omega$ . This idea is essentially contained in [5] in a different form. It was extended then to more general elliptic problems in  $\mathbb{R}^2$  in [6], to problems on a two-dimensional manifold in [9] and to the case of nonlinear boundary conditions in [10].

Let now  $X(s)$  be the solution of

$$(2.1) \quad \begin{cases} X''(s) = c^2 f(X) & \text{in } (0, s_0) \\ X'(0) = 0, X(s_0) = 1, \end{cases}$$

where at the moment the precise form of  $f(X)$  is not yet relevant.

As a first choice of a linear problem consider the “torsion problem”, i.e.

$$(2.2) \quad \begin{cases} \Delta\psi + 1 = 0 & \text{in } \Omega \\ \psi = 0 & \text{in } \partial\Omega. \end{cases}$$

One then constructs a supersolution  $\bar{u}(x)$  to (1.1) having the same level lines as the torsion function by setting

$$(2.3) \quad \bar{u}(x) = X(s(x)), \quad x \in \Omega,$$

where

$$(2.4) \quad s(x) = \sqrt{2(\psi_m - \psi(x))}, \quad \psi_m = \max_{\Omega} \psi.$$

The choice of  $s(x)$  is suggested by the one-dimensional version of (2.2) when  $s(x) = x$ . In problem (2.1) we thus choose  $s_0 = \sqrt{2\psi_m}$ .

The main result from which the estimates follow can be stated as

**Theorem 1** *Assume that the mean curvature of  $\partial\Omega$  is nonnegative everywhere and  $f(0) \geq 0$ ,  $f'(s) \geq 0$  for  $s \geq 0$ . Then*

$$\begin{aligned} \bar{u}(x) &= X(s(x)) \text{ is a supersolution, i.e.} \\ \Delta\bar{u} &\leq c^2 f(\bar{u}) \text{ in } \Omega \\ \bar{u} &= 1 \text{ on } \partial\Omega. \end{aligned}$$

*Proof:* Calculate first

$$\begin{aligned} \nabla s &= -\frac{\nabla\psi}{s}, \\ \Delta s &= -\frac{\Delta\psi}{s} - \frac{|\nabla\psi|^2}{s^3} = \frac{1}{s} - \frac{|\nabla\psi|^2}{s^3}, \end{aligned}$$

and then

$$\Delta\bar{u} = X' \cdot \Delta s + X'' \cdot |\nabla s|^2 = \frac{1}{s} X' \left(1 - \frac{|\nabla\psi|^2}{s^2}\right) + X'' \cdot \frac{|\nabla\psi|^2}{s^2},$$

from where one finds

$$\Delta\bar{u} - c^2 f(\bar{u}) = \left(\frac{1}{s} X' - c^2 f(X)\right) \left\{1 - \frac{|\nabla\psi|^2}{s^2}\right\}.$$

It was shown in [4] that under our assumption on  $\partial\Omega$  one has

$$|\nabla\psi|^2 \leq 2(\psi_m - \psi(x)) ,$$

which means that the term  $\{\}$  is nonnegative. It remains to check the sign of the other factor. To this end consider

$$g(s) = X' - c^2 s f(X)$$

which satisfies  $g(0) = 0$  and

$$g'(s) = X'(s) - c^2 f(X(s)) - c^2 \frac{df}{dX} \cdot X'(s) = -c^2 \frac{df}{dX} \cdot X'(s) \leq 0 ,$$

since  $\frac{df}{dX} \geq 0$  and  $X' \geq 0$ .

Therefore  $g(s) \leq 0$  for  $s \geq 0$  so that  $\bar{u}(x)$  satisfies the required differential inequality. If we select  $s_0 = \sqrt{2\psi_m}$  in (2.1) the boundary condition is satisfied as well.

For the particular choice  $f(u) = u^p$ ,  $0 \leq p < 1$  the usual properties of sub- or supersolutions still hold even if  $f'(0)$  becomes unbounded (see [11]).

As a first application of Theorem 1 we note

**Corollary 1.1** *The critical value satisfies*

$$(2.5) \quad c_0^2 \leq \frac{p+1}{(1-p)^2 \psi_m} ,$$

with equality if  $\Omega$  degenerates to an infinite slab of width  $2\sqrt{2\psi_m}$ .

*Proof:* For  $f(u) = u^p$  the value  $c_0$  in the one dimensional problem(2.1) is given by

$$(2.6) \quad c_0^2 = \frac{2(p+1)}{(1-p)^2} \cdot \frac{1}{s_0^2} .$$

Since  $\bar{u}(x) = X(s(x))$  is a supersolution inequality (2.6) follows immediately.

**Remarks:**

- a) In order to make inequality (2.5) more explicit we need a lower bound for  $\psi_m$ . It was shown in [4] that

$$\psi_m \geq \frac{A^2}{2L^2} \quad (A = |\Omega|, L = |\partial\Omega|)$$

with equality for a slab (taking appropriate limits) under our assumption on  $\partial\Omega$ . For a strictly convex plane domain it follows from inequality (3) of Webb [12] that even

$$\psi_m \geq \frac{A^2}{L^2} \cdot \frac{L - k_0 A}{2L - 3k_0 A}, \quad (\text{curvature } k \geq k_0 > 0)$$

with equality for a circle or an infinite strip.

b) It was shown in [3] that

$$(2.7) \quad c_0^2 \geq \frac{2(p+1)}{(p-1)^2 \cdot \rho^2}, \quad \rho = \text{radius of largest ball in } \Omega,$$

again with equality for a slab. Hence (2.5) is the optimal counterpart to (2.7) since  $\psi_m \leq \frac{c_0^2}{2}$  as noted in [4].

There is also information on the location and size of  $\Omega_0$  contained in Theorem 1, which may be stated as

**Corollary 1.2** *The dead core  $\Omega_0$  contains the set*

$$\left\{ x \in \Omega \mid \psi(x) \geq d(p, c) \left[ \sqrt{2\psi_m} - \frac{1}{2} d(p, c) \right] \right\},$$

$$\text{where } d(p, c) = \frac{\sqrt{2(p+1)}}{(1-p)c}.$$

*Proof:* In the one-dimensional problem (2.1) with  $f(X) = X^p$  one can easily calculate the dead core as the interval  $(0, \sigma(p, c))$  where

$$(2.8) \quad \sigma(p, c) = s_0 - d(p, c).$$

The level set  $\sqrt{2(\psi_m - \psi(x))} = s_0 - d(p, c) = \sqrt{2\psi_m} - d(p, c)$  must be contained in  $\Omega_0$  since  $X(s(x))$  is a supersolution. This implies the statement of Corollary 1.2.

**Remarks:**

a) It was shown in [3] that the dead core is contained in the set

$$\{x \in \Omega \mid \text{dist}(x, \partial\Omega) \geq d(p, c)\}.$$

b) The torsion function  $\psi(x)$  is only known explicitly in some special cases (e.g. ellipse, equilateral triangle). If  $\psi(x)$  is not known explicitly Corollary 1.2 is still useful if one makes use of the monotonic behavior of  $\psi(x)$  with respect to the domain: if  $\tilde{\Omega} \subset \Omega$  then the corresponding solutions satisfy  $\tilde{\psi}(x) \leq \psi(x)$  for any  $x \in \tilde{\Omega}$ .

As a third consequence of Theorem 1 one has

**Corollary 1.3** *For given value of  $\psi_m$  the effectiveness factor  $\eta$  is a minimum for the slab of width  $2\sqrt{2\psi_m}$ . In particular, if  $c \geq c_0$  one has*

$$(2.9) \quad \eta \geq \frac{1}{c\sqrt{(p+1)\psi_m}}.$$

*Proof:* Since  $\bar{u} = X(s(x))$  is a supersolution which satisfies the boundary condition one has

$$\frac{\partial u}{\partial n} \geq \frac{\partial \bar{u}}{\partial n} \quad \text{on } \partial\Omega$$

which implies

$$c^2|\Omega|\eta \geq - \oint_{\partial\Omega} \frac{X'(s_0)}{s_0} \frac{\partial \psi}{\partial n} d\sigma = \frac{X'(s_0)}{s_0} |\Omega|, \quad (d\sigma = \text{element of } \partial\Omega)$$

that is

$$\eta \geq \frac{X'(s_0)}{c^2 \cdot s_0} = \text{effectiveness factor for slab of width } 2 \cdot s_0.$$

If  $c \geq c_0$  one gets (see Aris [1], p. 146)

$$\eta = \sqrt{\frac{2}{c^2(p+1) \cdot s_0}} = \frac{1}{c\sqrt{(p+1)\psi_m}},$$

which completes the proof.

**Remarks:**

- a) For  $c < c_0$  the value of  $\eta$  in the one-dimensional case is determined from the relation (see [1], p. 148)

$$c \cdot s_0 = \sqrt{\frac{2}{1+p}} (1 - u_0^{p+1}) \cdot F\left(1, \frac{p}{p+1}; \frac{3}{2}; 1 - u_0^{p+1}\right) =: h(u_0, p)$$

and then,

$$\eta = \frac{1}{F\left(1, \frac{1}{p+1}; \frac{3}{2}; 1 - u_0^{p+1}\right)}$$

there  $F(a, b; c; z)$  is the hypergeometric function and  $u_0 = u(0) = \text{minimum value}$ . The function  $h(u_0, p)$  is monotonically decreasing in  $u_0$  for any  $p \in (0, 1)$ . One can also prove that  $\eta$  is a decreasing function of  $s_0$ , so that an upper bound for  $\psi_m$  is needed in Corollary 1.3. A number of upper bounds for  $\psi_m$  are known (see e.g. [2], [4], [8]). One has e.g.  $\psi_m \leq \frac{p^2}{2}$  under our assumptions on  $\partial\Omega$ .



b) It follows from Lemma 3.1 of [3] that for  $c \geq c_0$

$$(2.10) \quad \eta \leq \sqrt{\frac{2}{p+1}} \cdot \frac{L}{cA}$$

and the equality sign holds also in the limit if  $\Omega$  degenerates into a slab.

A second choice of a linear problem is the fixed membrane problem on  $\Omega$  i.e.

$$(2.11) \quad \begin{cases} \Delta\varphi + \lambda\varphi = 0 & \text{in } \Omega \\ \varphi = 0 & \text{on } \partial\Omega . \end{cases}$$

In this case we replace  $s(x)$  as defined in (2.4) by

$$(2.12) \quad t(x) = \frac{1}{\sqrt{\lambda_1}} \cos^{-1} \left( \frac{\varphi(x)}{\varphi_m} \right) ,$$

where  $\lambda_1 =$  first eigenvalue with associated eigenfunction  $\varphi(x)$  and  $\varphi_m = \max_{\Omega} \varphi$ . Also  $s_0 = \frac{\pi}{2\sqrt{\lambda_1}}$  is now the length of the interval.

The analogue of Theorem 1 is now

**Theorem 2** *Assume that the mean curvature of  $\partial\Omega$  is nonnegative everywhere and  $f(0) \geq 0$ ,  $f'(s) \geq 0$  for  $s \geq 0$ . Then*

$$\begin{aligned} \bar{u}(x) &= X(t(x)) && \text{is a supersolution, i.e.} \\ \Delta\bar{u} &\leq c^2 f(\bar{u}(x)) && \text{in } \Omega \\ \bar{u} &= 1 && \text{on } \partial\Omega \end{aligned}$$

*Proof:* A straightforward calculation gives now

$$\Delta\bar{u} - c^2 f(\bar{u}) = \left( \lambda_1 - \frac{|\nabla\varphi|^2}{\varphi_m^2 - \varphi^2} \right) \left\{ X' \cdot \cot(\sqrt{\lambda_1} t) \sqrt{\lambda_1} - c^2 f(X(t)) \right\} .$$

By a result of Payne & Stakgold [7] one has  $\lambda_1 \geq \frac{|\nabla\varphi|^2}{\varphi_m^2 - \varphi^2}$  if the mean curvature of  $\partial\Omega$  is nonnegative.

The term  $\{ \}$  is nonpositive as a similar reasoning as in the proof of Theorem 1 shows, now for  $g(t) = X' \sqrt{\lambda_1} \cos(\sqrt{\lambda_1} t) - \sin(\sqrt{\lambda_1} t) f(X(t))$ .

The counterparts of Corollaries 1.1 - 1.3 are now obvious:

**Corollary 2.1** *The critical value  $c_0$  satisfies*

$$(2.13) \quad c_0^2 \leq \frac{8\lambda_1(p+1)}{(1-p)^2\pi^2} .$$

**Remarks:** Since  $\lambda_1 \geq \frac{\pi^2}{8\psi_m}$  as noted by Payne [5], (2.13) is weaker than (2.5).

The counterpart of Corollary 1.2 may still be useful in some cases (see e.g. Example 1). It now reads

**Corollary 2.2** *The dead core is contained in the set*

$$\{x \in \Omega \mid \varphi(x) \geq \varphi_m \cdot \sin(\sqrt{\lambda_1} d(p, c))\} .$$

It is interesting to see however that inequality (2.9) can now be improved. In fact we have

**Corollary 2.3** *For given value of  $\lambda_1$  the effectiveness factor  $\eta$  is a minimum for the slab of width  $\frac{\pi}{\sqrt{\lambda_1}}$ . In particular if  $c \geq c_0$  one has*

$$(2.14) \quad \eta \geq \frac{2}{c\pi} \sqrt{\frac{2\lambda_1}{p+1}} .$$

### Examples

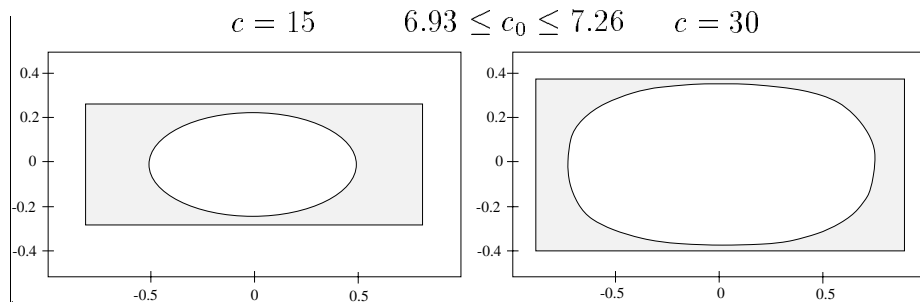
1.  $\Omega =$  Rectangle of sides  $a = 2$ ,  $b = 1$ . Inequalities (2.5) and (2.7) then yield

$$8 \leq \frac{(1-p)^2}{p+1} c_0^2 \leq 8.782 .$$

For  $c \geq c_0$  (2.10) and (2.14) give

$$\sqrt{\frac{2}{p+1}} \cdot \frac{2.236}{c} \leq \eta \leq \sqrt{\frac{2}{p+1}} \cdot \frac{3}{c} .$$

For  $p = \frac{1}{2}$  Corollary 1.2 and the estimate in [3] give the following pictures (the boundary of  $\Omega_0$  must lie in the shaded regions).



2.  $\Omega =$  Equilateral triangle of height 1. One has  $\psi_m = \frac{1}{27}$  and  $\lambda_1 = 4\pi^2$  and (2.5) and (2.7) now give

$$18 \leq \frac{(1-p)^2}{p+1} c_0^2 \leq 27 .$$

A better lower bound in some cases (for plane domains) is  $c_0^2 \geq \frac{4\pi}{(1-p) \cdot A}$  as given in Corollary 3.1 of [3].

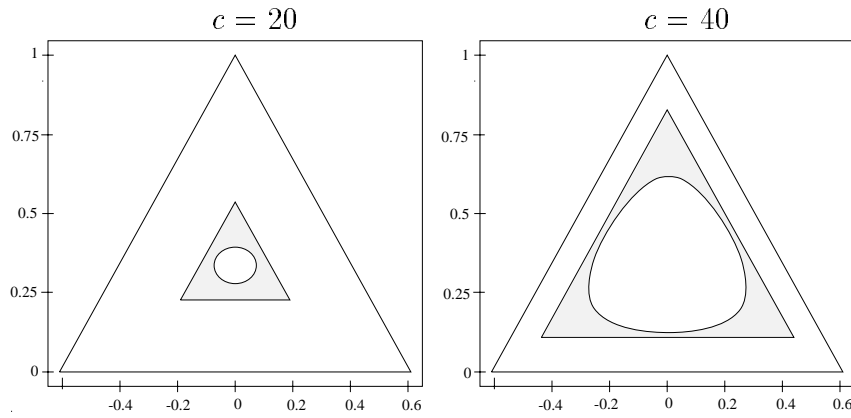
For  $c \geq c_0$  (2.10) and (2.14) show that

$$\sqrt{\frac{2}{p+1}} \cdot \frac{4}{c} \leq \eta \leq \sqrt{\frac{2}{p+1}} \cdot \frac{6}{c} .$$

For  $p = 0.5$  one has therefore

$$9.33 \leq c_0 \leq 12.73 ,$$

and one obtains the following pictures



The boundary of  $\Omega_0$  lies in the shaded region.

## References

- [1] Aris R., *The mathematical theory of diffusion and reaction in permeable catalysts*, Vol. 1, Clarendon Press, Oxford 1975.
- [2] Bandle C., *Isoperimetric inequalities and applications*, Pitman, London 1980.
- [3] Bandle C., Sperb R. and Stakgold I., *Diffusion and reaction with monotone kinetics*, *Nonlinear Analysis* 8, (1984), 321-333.
- [4] Payne L.E., *Bounds for the maximum stress in the Saint Venant torsion problem*, *Indian J. Mech. Math.*, Special Issue 1968, pp. 51-59.
- [5] Payne L.E., *Bounds for the solutions of a class of quasilinear elliptic eigenvalue problem in terms of the torsion function*, *Proc. Roy. Soc. Edinburgh* 1581, pp. 51-59.
- [6] Payne L.E. and Phillipin G., *Comparison theorems for a class of nonlinear elliptic problems*, *Nonlinear Analysis* 3, 1985, pp. 787-797.
- [7] Payne L.E. and Stakgold I., *Nonlinear problems in nuclear reactor analysis*, *Springer Lecture Notes in Math.*, Vol. 322, pp. 298-307.
- [8] Sperb R., *Maximum principles and their applications*, Academic Press, New York 1981.
- [9] Sperb R., *Optimal bounds for the critical value in a semilinear boundary value problem on a surface*, *General inequalities* 5, ISNM 80(1987), 391-400.
- [10] Sperb R., *Optimal bounds in semilinear elliptic problems with nonlinear boundary conditions*, *ZAMP* 44, 1993, pp. 639-653.
- [11] Stuart C.A., *Maximal and minimal solutions of elliptic differential equations with discontinuous nonlinearities*, *Math. Z.* 163 (1978), p. 239.
- [12] Webb J.R.L., *Maximum principles for functionals associated with the solution of semilinear elliptic boundary value problems*, *ZAMP* 40, 1989, pp. 330-338.